Liquid Phase Deposition of a Space-Durable, Antistatic $SnO₂$ Coating on Kapton

Katya Gotlib-Vainstein,† Irina Gouzman,‡ Olga Girshevitz,† Asaf Bolker,‡ Nurit Atar,‡ Eitan Grossman,‡ and Chaim N. Sukenik*,†

† Department of Chemistry a[nd](#page-6-0) Institute for Nanotechnology and Advanced Materials, Bar-Ilan University, Ramat Gan, Israel 52900 ‡ Space Environment Department, Soreq NRC, Yavne 81800, Israel

ABSTRACT: Polyimides are widely used in thermal blankets covering the external surfaces of spacecrafts due to their space durability and their thermooptical properties. However, they are susceptible to atomic oxygen (AO) erosion, the main hazard of low Earth orbit (LEO), and to electrical charging. This work demonstrates that liquid phase deposition (LPD) of 100 nm of tin oxide creates a protective coating on Kapton polyimide that has good adherence and is effective in preventing AO-induced surface erosion and in reducing electrical charging. The as-deposited tin oxide induces no significant changes in the original thermo-optical properties of the polymer and is effective in preventing electrostatic discharge (ESD). The durability of the oxide coating under AO attack was studied using oxygen RF plasma. The AO exposure did not result in any significant changes in surface morphology, thermo-optical, mechanical, and electrical properties of the tin oxide-coated Kapton. The erosion yield of tin oxide-coated Kapton was negligible after

exposure to 6.4 \times 10²⁰ O atoms·cm⁻² of LEO equivalent AO fluence, indicating a complete protection of Kapton by the LPD deposited coating. Moreover, the tin oxide coating is flexible enough so that its electrical conductivity stays within the desired range of antistatic materials despite mechanical manipulations. The advantages of liquid phase deposited oxides in terms of their not being line of site limited are well established. We now extend these advantages to coatings that reduce electrostatic discharge while still providing a high level of protection from AO erosion.

KEYWORDS: atomic oxygen, space environment, tin oxide, liquid phase deposition, electrostatic discharge

1. INTRODUCTION

Deployable structures and ultra-lightweight spacecraft require polymeric materials that possess a unique combination of mechanical and thermo-optical properties.¹ Polyimides like Kapton are very attractive for these applications because they are lightweight and flexible, with good wear [an](#page-6-0)d thermo-optical properties.² Kapton, which is used as an external layer of a multilayer isolation system, 3 is exposed extensively to the space environm[en](#page-6-0)t. The low Earth orbit environment (LEO) introduces different hazar[ds](#page-6-0) to external spacecraft materials, such as exposure to atomic oxygen (AO), UV and ionizing radiation, micrometeoroids, orbital debris, and charged particle bombardment.4−⁶ It is well-known that AO can cause significant erosion of Kapton;^{7,8} e.g., a 25 μ m thick Kapton blanket would [com](#page-6-0)pletely erode within 6 months in LEO.⁷ Thus, structures in the LEO [en](#page-6-0)vironment require protective coatings.

Various approaches have been developed to mitigate the AO erosion of polyimides. The protective strategies may be divided into three categories: (a) protective coatings, mainly metal oxides, produced by different methods; 7,9,10 (b) surface modification, such as ion implantation or direct application of siloxanes;^{11−13} (c) development of advanc[ed](#page-6-0) [hy](#page-6-0)brid materials that are stable in an oxidative environment, e.g., POSS/ polyimide nanocomposites.14−¹⁶

Avoiding or minimizing charging is an important consideration in the design of sp[acecra](#page-6-0)ft. Energetic charged particles can penetrate a near surface polymer layer, thus depositing a charge onto insulating materials. This can result in electrostatic charge (ESC) build-up, development of large electric fields, and eventually a discharge.^{17,18} Electrostatic discharge (ESD) can cause serious damage to delicate spacecraft electronics.¹⁹ The surface resistivity ne[eded](#page-6-0) to mitigate ESC build-up on insulators is in the range of 10⁶ to 10^{10} Ω \square ^{-1.20} The [w](#page-6-0)idely . used $SiO₂$ protective coating physically blocks the interaction with AO^{21} but does not prevent the ESC build[-u](#page-6-0)p due to its insulating properties. Indium tin oxide is a standard conductive coating; 2^{22} 2^{22} however, it is very brittle and loses its conductivity when folded to the point of cracking.²³ The combination of sufficie[nt](#page-6-0) electrical conductivity to diminish ESC build-up without affecting material flexibility is [di](#page-6-0)fficult to achieve.

Liquid phase deposition (LPD) is a commonly used method for preparing oxide ceramic films. It is technologically simple

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since it involves film deposition from aqueous solution under near-ambient conditions and allows coatings of substrates with complex shapes.²⁴ In previously published work, we demonstrated that a titania $(TiO₂)$ coating provides good protection of Kapton agai[ns](#page-6-0)t AO (mass loss about 1% compared to uncoated Kapton). 25 However, the cracking of the coating and its high resistivity renders a $TiO₂$ coating unsatisfactory for applications where [pr](#page-6-0)eventing ESD is critical. The present work demonstrates that LPD tin oxide $(SnO₂)$ overcomes both of these drawbacks.

In this work, we extend the previously reported LPD of $SnO₂$ films on substrates like silicon, glass, and gold electrodes,^{26–28} to coating Kapton. SnO_2 -coated Kapton is characterized in terms of its morphology, growth rate, chemical compo[sition,](#page-6-0) mechanical and thermo-optical properties, as well as AO durability. We also compare the electrical properties of the previously reported $TiO₂$ -coated Kapton with those of $SnO₂$ coated Kapton by Kelvin probe atomic force microscopy (KP-AFM), and electrical measurements. We also use Rutherford backscattering spectrometry (RBS) and the distortions in these spectra created by ion beam induced charging to compare variously treated Kapton films with each other and with simulated data so as to demonstrate the effectiveness of $SnO₂$ as a dissipative coating.

2. EXPERIMENTAL SECTION

2.1. Thin Film Preparation. TiO₂ and $SnO₂$ coatings were prepared on 125 μ m thick Kapton 500 HN sheets (DuPont) by LPD.^{24−26} Prior to coating, the as-received sheets were cut into 1 cm \times 1 cm squares or 1 cm \times 4 cm rectangular coupons, washed with doub[le](#page-6-0) [dis](#page-6-0)tilled water and ethanol, and dried under nitrogen. The coupons were pretreated by exposing them to an air plasma (Harrick, model PDC-3XG) at a pressure of 0.3 mmHg and 18 W power for 20 min.25,29,30 Immediately after plasma pretreatment, the samples were placed into the LPD solution at room temperature. The following solu[tions](#page-6-0) [w](#page-7-0)ere used

- (i) TiO₂: 0.3 M H₃BO₃ and 0.1 M $(NH_4)_2$ TiF₆ in water. Substrates were left in the solution for 8 h.
- (ii) SnO_2 : 0.03 M SnF_2 , 0.45 M H_3BO_3 , 0.15 M HF, and 0.06 M $H₂O₂$ in water. The deposition time varied from 1 to 6 h.

The coated samples were rinsed in water and methanol before drying under conditions of controlled humidity²⁹ for 53 h. The stability of the coating was confirmed by sonication in water for 10 min. The robustness and good adhesion of the coa[ting](#page-6-0) were tested. All the SnO₂−coated samples were found to be stable to sonication in water for 10 min and were sufficiently adherent so that they were not compromised by a standard tape test. 31

2.2. Characterization Techniques. 2.2.1. Scanning Electron Microscopy (SEM). The surface m[orp](#page-7-0)hology of the samples was assessed by SEM (Magellan, FEI), at an accelerating voltage in the range of 5−10 kV with a 10 nm surface gold coating.

2.2.2. X-ray Photoelectron Spectroscopy (XPS). Surface chemical composition was determined by XPS (Kratos AXIS-HS spectrometer) using a monochromatized Al K α source. All data acquisition was done in a hybrid mode (using electrostatic and magnetic lenses) and detection pass energies of 40−80 eV. All XPS measurements were carried out at room temperature, under pressure in the range of 1.0− 3.0×10^{-9} Torr. The spectra were acquired with an electron flood gun for charge neutralization. The spectrometer energy scale was calibrated according to the ISO TC/201 SC7 international procedure for binding energy (BE) calibration with Au $4f_{7/2}$ = 83.98 eV and Cu $2p_{3/2}$ = 932.67 eV. Data processing was done with VISION 2.1 software (Kratos) using sensitivity factors for quantification. In most cases a Shirley background was used.³² Curve fitting was performed using a 80/20 Gaussian/Lorentzian line shape.

2.2.3. Rutherford Backscattering Spectrometry (RBS). The thickness of the $SnO₂$ layers was measured by RBS, using $SnO₂$ atomic density. RBS was also used for comparing the charging of uncoated Kapton and Kapton coated with $TiO₂$ or $SnO₂$ films. This work was done using a 1.7 MV Pelletron accelerator (NEC, USA). All spectra were collected using a 2.023 MeV 4 He⁺ \pm 1 keV beam. The beam current was ∼10 nA, with a nominal beam diameter of 1.5 mm. An electron suppressor between the beam entrance and the sample holder was biased at 100 V vs ground. A second electron suppressor was placed in front of the sample, and it was biased at −300 V relative to the sample holder. RBS spectra were acquired using a fixed silicon drift detector (ULTRA Silicon-Charged Particle Detector, ORTEC) in a Cornell geometry with detector scattering angle 169° and solid angle of 2.7 msr. A normal incident beam was used for all measurements. All samples were mounted on the holder by double sided, self-adhesive carbon tape. In the case of film thickness studies, the charging effect on oxide-coated Kapton samples was compensated by a thin Au coating (8 nm) . Data Furnace computation code (NDFv9.6a program)³³ and SRIM 2003^{34} stopping powers were used to fit the data.

2.2.4. Atomic Force Microscopy (AFM). All scanning [pr](#page-7-0)obe microscopy [m](#page-7-0)easurements were performed using an ICON instrument (Bruker AXS SAS). The measurements were carried out in tapping, nanoindentation/scratching and surface potential modes. The rootmean-square roughness (R_q) was calculated from $1 \times 1 \mu m^2$ micrographs. The deflection sensitivity of each probe was measured by pressing the probe on a hard surface, and the spring constant was calibrated by the "Sader method".³⁵ Nanoscratching was done using a diamond coated tip (DDESP) (force constant of 20−80 N m[−]¹ , Digital Instruments, Santa Barb[ara,](#page-7-0) CA). Nanoscratching was done with the indenter at variable normal loads of $10-20$ μ N, a sliding speed of 0.3 μ m s⁻¹, and a scratch length of ~1.5 μ m. Nanoindentation was done with a Berkovich probe (force constant of 194 N m⁻¹) with applied load of 5 μ N. The stiffness value was calculated based on 18 points of indentation for each sample. The same indenter was used to image the area after the nanomechanical tests. The electrical measurements were done using scanning Kelvin Probe AFM under ambient conditions. The surface potential was determined using conductive probes coated with cobalt/chromium (MESP; force constant of 1−5Nm[−]¹ , Digital Instruments, Santa Barbara, CA). The lift scan height was adjusted to 50 nm, and the applied voltage on the tip was 1 V. Before analysis of the images, first order "flatten" and "planefit" functions were applied to each image. The roughness was determined by Nanoscope analysis software.

2.2.5. Thermo-Optical Measurements. Solar absorptance (α_s) and thermal emittance (ε) were measured using a TASA 2000 portable reflectometer (AZ Technology, Inc.).

2.2.6. Surface Resistivity. Macroscopic electrical properties of the samples in the lateral direction were measured using a source meter (Keithly 2400), under ambient conditions. In order to determine sheet resistivity, a set of parallel electrodes was applied across the surface of a 0.5×4 cm² sample to form several square segments with area of about 25 mm² each. The sheet resistance, R, was measured for different areas with a constant width (5 mm) , w, and increasing length, L (up to 4 cm). The sheet resistivity, ρ_s , in units of $\Omega \square^{-1}$ (ohm per a single square segment) is related to the sheet resistance according to the relationship:

$$
R = \rho_s \frac{L}{w} \tag{1}
$$

 ρ_{s} was calculated as the slope of the linear trend line that was fitted to the resistance versus the number of measured segments, L/w , and the contact resistance is expressed by the intercept of the linear line with the ν axis.

2.3. Atomic Oxygen Exposure. The AO exposure facility is based on a LB-3000 Advanced Energy RF-plasma system with a feed gas of 99.999% pure oxygen. The system was operated at a pressure of 40 mTorr, power of 1000 W, and oxygen flow of 10 sccm. Redirection of the afterglow through right angle deflection results in a strong reduction of ion current and UV radiation flux, facilitated by a supply of electrons from the metallic chamber walls, and radiation absorption

by the walls, respectively. The afterglow was characterized by optical emission spectroscopy, electrical measurements, UV radiation measurements, and Kapton etching rate measurements. A detailed description of the AO simulation system is presented elsewhere.³⁶ In the present work, the samples were exposed using two exposure regimes. In the first regime (regime I), special care was taken to e[xpo](#page-7-0)se the samples predominantly to AO atoms, without electron and ion fluxes originated from RF plasma. The LEO equivalent AO flux in this case is low, about 5 \times 10¹³ atoms s⁻¹cm⁻². Relatively low LEO equivalent AO fluence (about 4.3×10^{19} atoms cm⁻²) was achieved using this regime, and it was used for continuous online monitoring of mass losses from uncoated and $SnO₂$ -coated samples. To increase AO flux, the experiment was carried out under severe exposure conditions (regime II), using a direct RF $O₂$ plasma afterglow. The LEO equivalent AO flux in this experiment was 6.18 \times 10¹³ atoms s⁻¹ cm⁻² and the total AO fluence was 6.4×10^{20} atoms cm⁻². .

Atomic oxygen fluence measurements were conducted based on Kapton-HN mass loss, assuming an erosion yield of 2.80 \times 10^{−24} cm³ O atom⁻¹ and Kapton HN density of 1.42 g cm^{-3,37} The erosion yield . was determined gravimetrically, using an analytical balance (Sartorius SE2) with an accuracy of ± 0.1 μ g. All the measur[em](#page-7-0)ents were carried out immediately after sample removal from the exposure facility. The time of sample exposure to ambient air did not exceed 2 min.

For measuring the kinetics of mass loss, polyimide films were spincoated onto quartz crystal microbalance (QCM) crystals using a Dupont procedure for deposition of polyimide (Pyralin PI 2545, HD MicroSystems; $4-5 \mu m$ thick).³⁸ The deposited polyimide films were shown to be similar to Kapton HN films based on their FTIR sp[e](#page-7-0)ctra.³⁹ SnO₂ coatings were deposited on Pyralin-coated QCM crystals, using the same protocol as had been applied to the Kapton sheets.

3. RESULTS AND DISCUSSION

3.1. $SnO₂$ Coated Kapton. 3.1.1. Growth Rate and Surface Morphology. Commercial Kapton was coated with uniform, adherent, amorphous $SnO₂$ by LPD. Air plasma pretreatment of the Kapton samples was needed to ensure good adhesion of the SnO₂ to the Kapton. The $\left[\text{SnF}_{6}\right]^{-2}$ complex that is necessary for LPD was prepared in situ by reacting $SnF₂$ with H_2O_2 in the presence of HF. An overall ratio of $F/Sn = 7$ in solution is favorable for creating a uniform, crack-free (by HRSEM), oxide film on the Kapton surface. Figure 1 shows HRSEM and AFM images of the $SnO₂$ coating on Kapton obtained after 6 h of LPD. Both the HR-SEM and the AFM confirm that we obtain good surface coverage. Surface roughness, estimated by AFM, is about $R_q = 4.7$ nm, a little over 3 nm rougher than the original plasma treated Kapton (1.5 nm). This additional roughness is negligible relative to the overall thickness of the deposited oxide films.

The thickness of $SnO₂$ films as a function of the deposition time was evaluated by RBS. Figure 2a shows RBS spectra (only the tin region is shown) measured for SnO_2 -coated Kapton after different deposition times. [Th](#page-3-0)e width of the RBS peak (full width of half-maximum) was used to calculate the film thickness.⁴⁰ A minimum deposition time of about 2 h was required to create a fully covered surface. Increasing deposition time in t[he](#page-7-0) 2–6 h range results in a linear growth of the $SnO₂$ thickness. The short deposition time (0−2 h) deviated from the linear behavior (Figure 2b).

The slower growth rate at the beginning of the deposition process is consistent wi[th](#page-3-0) a film growth mechanism involving solution phase nucleation. This has been previously reported for titania deposition.⁴¹ Such a process is easy to control and provides predictable film thicknesses.

3.1.2. Atomic Ox[yg](#page-7-0)en Durability. $SnO₂$ deposited on a Pyralin-coated QCM crystal was subjected to AO exposure

Figure 1. Representative HR-SEM (a) and AFM (b) images of $SnO₂$ coatings (after 6 h coating time) on Kapton.

(regime 1). The total AO fluence was 4.3×10^{19} atoms cm⁻². . This fluence is equivalent to exposure of a satellite external surface to the LEO environment for about 2 months in an orbit of 500 km in the ram direction. Figure 3 shows mass loss measurements of two samples with different thicknesses of $SnO₂$ coatings (70 and 100 nm) as a [f](#page-3-0)unction of LEO equivalent AO fluence. Both $SnO₂$ films, 70 and 100 nm, provide barrier layers that protect polyimide from AO. However, the 70 nm film, despite providing a fully covered surface, showed measurable mass loss (15% of unprotected Pyralin after exposure to a fluence of 4.3×10^{19} atoms cm⁻²). Thicker, 100 nm, $SnO₂$ films provided complete protection against AO (Figure 3).

In a different experiment, the erosion of $SnO₂$ -coated Kapton was measured gravi[me](#page-3-0)trically after exposure to a higher fluence of 6.4 \times 10²⁰ atoms cm⁻² using the severe exposure conditions (regime II). This fluence is equivalent to exposure of a satellite external surface to the LEO environment for about 2 years in an orbit of 500 km in the ram direction. Even under these conditions, practically no erosion of tin oxide-coated (100 nm) Kapton was observed; the erosion yield of $SnO₂$ -coated Kapton is 0.3% of that measured for uncoated Kapton. According to our previous work, $42,43$ we can confidently assume that SnO_2 coated Kapton will be durable under exposure to much higher AO fluences (in [the](#page-7-0) range of 10^{21} atoms cm⁻²) in LEO environment, or under exposure to hyperthermal AO using a Laser detonation source.⁴⁴

The surface morphology of the AO exposed $SnO₂$ -coated samples was measured [b](#page-7-0)y electron microscopy. HR-SEM showed that as deposited 70 and 100 nm $SnO₂$ coatings on Kapton are uniform and crack-free (as seen in Figure 1a). While the thicker coating remained crack-free and showed no morphological changes after AO exposure, the 70 nm $SnO₂$ film

Figure 2. (a) $SnO₂$ peak in the RBS at various deposition times; (b) $SnO₂$ film thickness as a function of the deposition time, derived from RBS measurements.

Figure 3. Pyralin and $SnO₂$ -coated (70 and 100 nm) Pyralin mass losses as a function of AO fluence.

showed minor cracking (<1%). Those cracks (approximately 50 nm wide) are likely responsible, at least in part, for the observed mass loss. No prominent undercutting of Kapton was observed in the cracked regions; in contrast to the substantial cracking of TiO₂-coating and undercutting of Kapton that was reported previously.25,45 The seemingly lower protection efficiency of the thinner coating (70 nm) might be due, at least in part, to diffusion of AO through the oxide layer. Further experimentation would be needed to precisely define the optimal thickness of such coatings and to sort out the details of the factors that may contribute to the need for a minimal thickness. Nevertheless, the observation that 100 nm provides good protection is secure.

The 100 nm SnO_2 -coated Kapton, before and after AO exposure, was studied by XPS. The quantitative analysis of surface composition was carried out using the detailed XPS scans of C1s, O1s, and $Sn3d_{5/2}$ and standard atomic photoionization cross-section values. The results are summarized in Table 1. As-deposited $SnO₂$ coating showed a

Table 1. Surface Elemental Composition (at. %) Based on XPS Analysis

		Sn		O/Sn
Kapton + $SnO2$	41.3	19.4	39.4	
Kapton + $SnO2$ after AO	43.5	20.2	36.3	

substantial concentration of carbon on the surface. However, the absence of N1s signal indicated that this carbon originated from adventitious carbon contamination of the coated surface rather than from the underlying Kapton substrate. Based on the position and the line shapes of the O1s and Sn3d peaks, it is concluded that no chemical changes occur in the $SnO₂$ coating after AO exposure. In both cases, the peak at 487.5 eV corresponds to Sn^{4+} and the peak at 531.5 eV corresponds to O^{2−}. The ratio of oxygen-to-tin showed minimal change, from 2.1 to 2.2. A decrease in the carbon content, and a slight increase in oxygen may be explained by losses of adventitious carbon and postexposure water adsorption, respectively.

Having ascertained that the 100 nm thickness was essential for good AO protection, all subsequent characterizations were done with samples of this thickness. This means that the mechanical, thermo-optical, and electrical measurements reported below were all done on 100 nm thick samples before and after exposure to 4.3×10^{19} atoms cm⁻² AO fluence.

3.1.3. Mechanical Properties. The $SnO₂$ -coated samples were scratch tested before and after AO exposure. Figure 4 shows AFM images of $SnO₂$ -coated Kapton samples after scratching with forc[es](#page-4-0) of 10, 15, and 20 μ N. Both samples showed visible scratch grooves at all the applied loads. However, the measured groove depth of the scratches showed some differences. The 10 μ N scratch is very ill defined on the as-deposited $SnO₂$ with scratch depths of 10.4 and 27.1 nm for 15 and 20 μ N applied loads. The AO exposed SnO₂ showed grooves of 10.3, 15.5 and 29.7 nm, for 10, 15, and 20 μ N, respectively. The AO exposed samples showed slightly deeper scratch grooves. Such small changes do not suggest that the SnO2 undergoes serious mechanical changes after AO exposure. Those changes are in the range of the roughness of the $SnO₂$ film $(R_a = 4.7$ nm). Note that SnO₂ coating improves the scratch resistance of pristine Kapton which can be scratched even with 5 μ N (data not shown).²⁵ Thus, while these results suggest that this improved scratch resistance due to the $SnO₂$ coating may be somewhat under[min](#page-6-0)ed by the AO exposure, further work is needed to clarify this effect and its durability.

3.1.4. Thermo-Optical Properties. Thermo-optical properties, namely, solar absorptance (α_s) and thermal emittance (ε) , are responsible for the passive thermal control of spacecraft components. The equilibrium temperature, T (K), of SnO₂-

Figure 4. AFM images of scratch grooves and their depth on $SnO₂$ coated Kapton, before (a) and after (b) AO exposure.

coated Kapton in space is governed by the heat it absorbs from the sun and the heat it emits to its surroundings. It is calculated bv^4

$$
T = \left(\frac{\alpha_s}{\varepsilon}\right)^{1/4} \left(\frac{SA_n}{\sigma A}\right)^{1/4} \tag{2}
$$

where α_s is related to the film's solar absorptance, ε is its emittance, S (W m[−]²) is the solar flux per unit area at the spacecraft orbit, A_n (m^2) is the film's surface area normal to the solar flux, A (m^2) is its total surface area, and σ (5.67 \times 10⁻⁸ W m^{-2} K⁻⁴) is the Stefan–Boltzmann constant. The calculated maximum temperature is related to a normal position and solar flux, S, of 1366.1 W m⁻², the average value above the earth's atmosphere at 1 AU distance.⁴⁶

Thermo-optical properties of Kapton and $SnO₂$ -coated Kapton, before and after A[O](#page-7-0) exposure, are summarized in Table 2. SnO₂ coating affects the α_s/ε ratio, so that the value $\alpha_s/\varepsilon = 0.507$ (for uncoated Kapton) is increased to $\alpha_s/\varepsilon =$ 0.763. Subsequently, the equilibrium temperature of the $SnO₂$ -

Table 2. Thermo-Optical Properties and the Equilibrium Temperature, $T(K)$, in Space of Kapton and SnO₂-Coated Kapton, before and after AO Exposure

sample	$\alpha_{\rm s}$	ε	$\alpha_{\rm s}/\varepsilon$	T(K)
Kapton	0.426	0.840	0.507	332.4
Kapton/AO	0.538	0.850	0.634	351.6
Kapton/SnO ₂	0.645	0.845	0.763	368.2
Kapton/SnO ₂ /AO	0.66	0.848	0.778	370.0

coated Kapton is increased by 35.8 °C, a reasonable temperature for a thermal blanket. The increase of the α_s/ε might be related to the slightly increased roughness of the oxide coating. After AO exposure, the α_s/ε ratio for uncoated Kapton increases from 0.507 to 0.634, while the coated sample does not show a significant change ($\alpha_s/\varepsilon = 0.778$). Stable thermo-optical properties are consistent with the AO durability of the coating.

3.1.5. Surface Resistivity. The surface resistivity of the $SnO₂$ coated Kapton was measured before and after AO exposure as described above (section 2.2.6). The sheet resistivity of uncoated Kapton is about 10^{16} Ω \Box^{-1} .⁴⁷ SnO₂ coated Kapton . showed a decrease of 8 orders of magnitude in surface resistivity ($\rho_s = 2 \times 10^8 \Omega \square^{-1}$). This [va](#page-7-0)lue slightly increased after AO exposure $(\rho_s = 4 \times 10^8 \,\Omega \square^{-1})$; however, it stays well inside the limits for antistatic materials.¹⁸

This result is an important factor in recommending the potential use of $SnO₂$ as a barrier coa[tin](#page-6-0)g against AO. While both $TiO₂$ and $SnO₂$ have comparable capabilities in preventing AO erosion of the underlying polyimide, $SnO₂$ also provides a meaningful decrease in surface resistivity, thus making it an antistatic coating that can protect against ESD.

3.2. Comparison between $TiO₂$ and $SnO₂$ Coatings on Kapton. 3.2.1. Kelvin Probe AFM Measurements. In addition to the above resistivity measurements, the electrical properties of the $TiO₂$ and $SnO₂$ -coated Kapton were assessed using Kelvin probe measurements.⁴⁸ To evaluate the surface potential, half of each sample was coated with a thin (∼5 nm) gold layer by sputtering; t[he](#page-7-0) Au coated side was used as a reference for the KP-AFM measurements. The potential images were recorded using conductive tips and a lift mode with lift height of 50 nm. The potential differences between Au and $TiO₂$ and between Au and $SnO₂$ are shown in Figure 5, panels a and b, respectively. It is important to note the different scales of the y-axes of these plots.

The potential diffe[re](#page-5-0)nce between Au and $TiO₂$ is greater than that measured between Au and $SnO₂$ (61.8 vs 26.8 mV). The higher surface potential indicates a greater extent of surface charging. Those results speak to the better antistatic properties of the $SnO₂$ coating. This conclusion is in good agreement with previously reported electrical properties of TiO₂^{49,50} and SnO₂ $\lim_{s \to 5}$ films.^{51,52} It is also consistent with the electrical measurements reported above.

3.[2.2. R](#page-7-0)BS Assessment of Surface Charging. RBS typically provides a powerful analysis method for determining surface atomic composition and elemental depth profiles.⁴⁰ When bombarding a nonconductive sample with high energy ions, there is charge accumulation on the surface. The RB[S a](#page-7-0)nalysis of insulating materials is complicated by surface charging and spectral distortion.^{53–55} A correct simulation of distorted spectra is impossible; however, we demonstrate below how we make use of sp[ect](#page-7-0)r[al](#page-7-0) distortion in order to compare charge build up on a sample surface. In this way, we are not only measuring electrical properties (as described above) but are directly assessing how well a material mitigates charge build-up.

The energy of the detected backscattered particles in RBS is given by the equation

$$
E_{\rm d} = kE_0 \tag{3}
$$

where k is the kinematic factor and E_0 is the incident energy. Due to the build-up potential φ , the incident positive He ions of charge e undergo a deceleration and strike the sample surface with the energy $(E_0 - e\varphi)$. On the other hand, the backscattered particles are accelerated by the surface potential

Figure 5. KP-AFM potential images of Au-TiO₂ coating (a) and Au- $SnO₂ coating (b).$

ne φ , where *n* is the He charge state +1 or +2.⁵⁶ This results in a higher detected energy for the backscattered particles from the surface.^{55,57}

$$
E_{\rm d1} = k \times (E_0 - e\varphi) + n e\varphi \tag{4}
$$

Therefore, the energetic shift of experimental data represents the build-up potential due to a charge accumulating effect. The surface potential of the samples can be measured as the difference between the experimental data and the simulated data, as can be seen in Figure 6.

Uncoated and oxide coated Kapton samples were irradiated with a 2.0 MeV ⁴He⁺ ion beam, and the surface potential was measured as the difference between simulated and experimental data. The surface potential of Kapton was 27.5 kV. The $TiO₂$ coated sample had a slightly different potential, 25.1 kV . SnO₂coating of similar thickness decreases this potential to 18.3 kV (Table 3). The lower value of surface potential observed for the $SnO₂$ -coated sample indicates that it can be considered as coating capable of preventing electrostatic discharge.

3.3.3. Elastic Properties of TiO₂ and SnO₂ Coated Kapton. It is important to preserve the elastic properties of the Kapton after depositing the protective layer in order to allow handling of the material without degrading its electrical properties. The elastic measurements were performed on $TiO₂$ and $SnO₂$ coated Kapton by AFM nanoindentation. The stiffness of both uncoated Kapton and oxide-coated Kapton samples was estimated from nanoindentation force−displacement curves.

Figure 6. 2 MeV ⁴He RBS spectra and simulation of pristine Kapton. Built-up potential φ was estimated by the energetic shift between experimental and simulated data.

Table 3. Surface Potential

The slope of the unloading curve is the stiffness of the sample being measured.⁵⁸ The applied load was 5 μ N to avoid the substrate influence on the coated samples. The Kapton stiffness was measured t[o b](#page-7-0)e 201 \pm 15 N m⁻¹, the stiffness of TiO₂coated Kapton was 467 \pm 47 N m⁻¹, and the stiffness of SnO₂coated Kapton was 347 ± 24 N m⁻¹. Those results indicate that the $SnO₂$ coating is more flexible, compared to a $TiO₂$ coating, thus making it a better candidate for use as a protective layer.

Furthermore, the SnO_2 -coated Kapton was tested to explore the effects of winding around a rod on surface resistivity. Table 4 shows the resistivity of the coated Kapton as a function of

Table 4. Effects of Windings around the Ceramic Rod on Surface Resistivity

sample	Kapton/ $SnO2$ as	Kapton/ $SnO2$ after	Kapton/ $SnO2$ after
	deposited	5 windings	10 windings
resistivity $(\mathrm{G}\Omega\ \Box^{-1})$	0.2	3.5	21.7

windings around a ceramic rod of 2.9 mm diameter. The resistivity of the SnO_2 -coated Kapton increases after 5 and 10 windings around the rod, probably due to minor migration of the oxide grains; however it did not get to the value characteristic for uncoated Kapton $(10^{16} \Omega \square^{-1})$.⁴⁷ This means that the coating retains its antistatic characteristics after 10 windings around the ceramic rod. $59,60$

4. CONCLUSIONS

Uniform, adherent, crack-free $SnO₂$ coatings with thicknesses approaching 100 nm were deposited on Kapton. The coating was characterized in terms of surface morphology, mechanical and thermo-optical properties, AO exposure durability, and chemical composition. The 100 nm thick coating was found to provide nearly complete protection of Kapton under exposure to LEO equivalent AO fluence of 6.4 \times 10²⁰ atoms cm⁻², with

an erosion yield of only 0.3% of that measured for unprotected Kapton. No changes in chemical composition, surface morphology, scratch resistance, or thermo-optical properties were observed after AO exposure. The surface resistivity of SnO2-coated Kapton was found to be 8 orders of magnitude lower than that of uncoated Kapton (0.2 $\rm G\Omega$ \Box^{-1}), making it an antistatic material. The antistatic properties of $SnO₂$ -coated Kapton were preserved even after AO exposure and after mechanical manipulations.

Charge dissipation properties of $SnO₂$ -coated Kapton were compared to TiO₂-coated and uncoated Kapton films using Kelvin probe AFM and RBS. The surface potential of a $SnO₂$ coating was half that measured for a $TiO₂$ -coating. Charging effects in RBS analyses were used to assess charge build up, and they confirm that there is less charge build up on $SnO₂$ -coated Kapton compared with a $TiO₂$ coating or with uncoated Kapton.

The importance of our findings is that LPD $SnO₂$ can be successfully deposited on Kapton with good adhesion and that it is able to provide barrier properties that are as good or better than $TiO₂$ while significantly lessening the problem of ESD. Such coatings have great promise for space materials applications.

■ AUTHOR INFORMATION

Corresponding Author

*E-mail: chaim.sukenik@biu.ac.il.

Notes

The auth[ors declare no competin](mailto:chaim.sukenik@biu.ac.il)g financial interest.

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